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# Investigation of Role of Tungsten Ions on Structural and Optical Properties of Sodium Borosilicate Germinate Glass for Optoelectronic Applications

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#### Abstract

In the current work, the role of tungsten ions on the structural and optical properties of sodium borosilicate germinate glass with the composition [70 Na <sub>2</sub>B<sub>4</sub>O<sub>7</sub>–15 SiO<sub>2</sub>–(15-x) (Ge<sub>2</sub>O<sub>3</sub>) – x (WO<sub>3</sub>) while, x = 0,2,4, 6, 8 mol %] were studied. Fast quenching method were used to prepare the glass samples. Experimental and empirical density results confirm the amorphous nature of the prepared samples. Fourier transform infrared, FTIR, results showed  $N_4$  decreases as  $WO_3$ increases. These results suggest that the decreasing in non-bridging oxygen (NBO), back conversion  $BO_4$  to  $BO_3$ , occur by the increase of WO<sub>3</sub>. Optical band gap show decreases from 3.3 down to 1.89 eV with increase in refractive index from 1.98 up to 2.35 these results indicate promising application for the prepared glasses in electric and optoelectronic applications.

Keywords: Optical properties, Borosilicate glass, Fourier transform infrared

### 1. Introduction

G lass is an amorphous, transparent, brittle and<br>colourless material. Also, the structural properties of the glass can be improved by altering some components during the manufacturing process. Glass has versatile optical properties which is ideal for different optical applications. Depending on the purity and composition, glass can transmit light in both Infra-red and visible region of spectra. Borosilicate glasses have higher crystallization resistance compared with other glasses, which is useful in various technological applica-tions [\[1\]](#page-9-0). Na<sup>+</sup> ion can break down the original network of boron by acting as charge compensator in the four coordinated boron  $[2-4]$  $[2-4]$  $[2-4]$ . Transition metal oxides in glass have different important applications  $[3-5]$  $[3-5]$  $[3-5]$ . While borate glasses have low melting point and high mechanical strength  $[6-11]$  $[6-11]$  $[6-11]$ . Bismuth ions in borate glasses have high

density which nominate these glasses to be used as radiation shielding  $[12-14]$  $[12-14]$  $[12-14]$  $[12-14]$ . Tungsten oxide as a transition metal oxide has different valence states that play a important role in the electronic conduction. While Na<sup>+</sup> ions side by side with the  $BO_4$ and nonbridging oxygens increase ionic conductions [\[15](#page-9-5)]. In this paper, the effect of tungsten ions on the optical absorption, density, infrared spectra, and optical energy of tungsten borosilicate germinate glass system was reported. The electrical properties of the obtained system, which is already published [[15\]](#page-9-5), along with the optical properties obtained in the current work, the prepared glass system could be very useful in many optical and optoelectronic application.

#### 2. Experimental work

Glass samples with the composition [70  $Na_2B_4O_7-15 SiO_2 - (15-x) Ge_2O_3 - (X) WO_3$ ] where,  $x = 0$ , 2, 4, 6, and 8 mol. % were prepared by the

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melt quenching technique. High purity of row oxides of  $Na<sub>2</sub>B<sub>4</sub>O<sub>7</sub>$ ,  $Ge<sub>2</sub>O<sub>3</sub>$ ,  $SiO<sub>2</sub>$ , and  $WO<sub>3</sub>$  were used. Samples were melted in porcelain crucible at 1200  $^{\circ} \mathsf{C}$  for one hour. The molten was poured and quenched between two brass plates. Experimental density of the samples was measured by Archimedes's method using Toluene as an immersion liquid (FTIR) absorption spectra were obtained, by KBr method, in the wavenumber range of  $400-2000$  cm<sup>-1</sup> using Bruker vertex80 v. The optical absorption and transmission spectra were measured at wavelengths between 200 and 1100 nm for the bulk glasses using Shimadzu 3600 UV $-V$ is spectrophotometer. The appearance of sample  $x = 0$  is shown below.



#### 3. Results and discussions

#### 3.1. Density and molar volume

The density of the examined glasses was evaluated empirically using the liquid displacement method, and the results were compared to those found theoretically for the compositions' close-packed structures. For comparison, the experimental and empirical density values were both displayed in [Fig. 1](#page-2-0) as a function of  $WO<sub>3</sub>$  content. It is evident that both density values experimental and empirical grew steadily and linearly as  $WO<sub>3</sub>$  concentration increase for the compositions' close-packed structures. For comparison, the experimental and empirical density values were both displayed in [Fig. 1](#page-2-0) as a function of  $WO<sub>3</sub>$  content. It is evident that both density values experimental and empirical grew steadily and linearly as  $WO<sub>3</sub>$  concentration increase [\[16](#page-9-6)]. It is also noted that the empirical density is typically higher than the corresponding experimental density, as shown in [Fig. 1](#page-2-0) which confirm the short-range order in the samples under investigation. The difference between empirical and experimental density be explained as, the empirical density uses the crystalline density values of the glass component oxides while glass contains disorder with increase the overall volume which decrease the density of the glass than the empirical values. The behaviour of the density is expected with the addition of  $WO_3$  content due to the replacement of the lighter weight Ge ions Appearance of sample  $X = 0$ . with the heavier W ions.

<span id="page-2-0"></span>

Fig. 1. The change in density as function of  $WO<sub>3</sub>$  content.

Molar volume is closely connected to the internal spatial structure of materials, so, it is appropriate to provide the molar volume values of the examined glasses as well [[16\]](#page-9-6). The computed values of the molar volumes for both empirical and experimental calculations are shown in [Fig. 2](#page-3-0). As the  $WO_3$  content increases, the molar volume of empirical calculation exhibits a linear behaviour, while the experimental measurements show a progressive drop. Molar volume may be reduced because of the formation of the shorter atomic links, which may have increased the stretching force constant of the glass network bonds [\[15](#page-9-5)].

The oxygen packing density (O P D) values can be obtained by the following equation [\[17](#page-9-7)],

$$
O.P.D = \frac{\rho.O}{M_i} \tag{1}
$$

where, O is the number of oxygens per unit Formula.

The obtained values of O.P.D, given in Table (1), show a decreasing trend with the gradual addition of  $WO_3$ .

<span id="page-3-1"></span>The values of the average boron-boron separation distance  $(d_{B-B})$  were also obtained by using equation [\(2\),](#page-3-1)

<span id="page-3-0"></span>
$$
d_{B-B} = \left[\frac{V_m^b}{N_A}\right]^{\frac{1}{3}}\tag{2}
$$

where  $N_A$  is Avogadro's number and  $V_m^b$  is the volume of boron per mole of the sample, which it can be calculated using equation [\(3\)](#page-3-2),

<span id="page-3-2"></span>
$$
V_m^b = \frac{V_m}{2(1 - \chi_B)}\tag{3}
$$

where  $X_B$  is the mole fraction of boron [[16\]](#page-9-6).

The  $W^{3+}$  ion concentration in all samples was also calculated using equation [\(4\),](#page-3-3)

<span id="page-3-3"></span>
$$
N_W = n \left[ \frac{wt_{WO_3}}{M w_{WO_3}} \right] \rho N_A \tag{4}
$$

where  $Wt_{WO_3}$  is the percentage weight in the glass sample and  $Mw_{WO_3}$  is the molecular weight of WO<sub>3</sub>.

<span id="page-3-4"></span>The Polaron radius  $(r_p)$  can be calculated from equation [\(5\),](#page-3-4)

$$
r_p = \frac{1}{2} \left(\frac{\pi}{6N}\right)^{\frac{1}{3}}\tag{5}
$$

where N is the number of  $W^{3+}$  ion concentration in the sample [[18\]](#page-9-8).

The inter-nuclear distance is given by,

$$
r_i = \left(\frac{1}{N}\right)^{\frac{1}{3}}\tag{6}
$$

while the field strength (F) can be calculated as follows,



Fig. 2. The change in molar volume as a function of  $WO<sub>3</sub>$  content.

<span id="page-4-0"></span>Table 1. The obtained values of the measured and calculated physical parameters of the glasses under study.

Physical parameter for	S1	S <sub>2</sub>	S3	S4	S5
the glass samples.					
$\rho_{exp}$ g/cm <sup>3</sup>	2.57	2.60	2.64	2.66	2.69
$\rho_{\rm emp}$ g/cm <sup>3</sup>	3.02	3.03	3.03	3.04	3.04
$(V_m)_{emp}$ cm <sup>3</sup> /mole	59.17	59.16	59.15	59.15	59.14
$(V_m)_{exp}$ cm <sup>3</sup> /mole	69.54	68.94	67.97	67.53	66.77
$d_{B-B}$ (10 <sup>-8</sup> m)	4.77	4.75	4.73	4.72	4.70
N. of W $(10^{20}$ cm <sup>-1</sup> )		1.26	1.27	1.28	1.30
$r_p(10^{-7} m)$		2.00	1.99	1.98	1.98
$r_i$ (10 <sup>-8</sup> m)		8.05	8.01	7.99	7.96
F $(10^{16}$ cm <sup>-2</sup> )		1.14	1.15	1.16	1.17
OPD $\times$ 10 <sup>22</sup> (Atom/cm <sup>3</sup> )	8.66	8.73	8.86	8.91	9.02

$$
F = \frac{Z}{r_p^2} \tag{7}
$$

where z in the atomic number.

[Table 1](#page-4-0) shows the field strength and polaron radius of the glasses under investigation. It is noted the field strength increases with the increase of  $WO<sub>3</sub>$ ratio. This increase is due to the increase in interatomic distance and to the increase in the number of NBO ions because of the convert of  $BO<sub>4</sub>$  to  $BO<sub>3</sub>$ .

#### 3.2. FTIR spectra

<span id="page-4-1"></span>[Figure 3](#page-4-1) displays the IR absorption spectra for glass samples spanning the 400 to 2000  $cm^{-1}$ wavenumber range. The sample  $S_1$  deconvoluted FTIR spectrum, seen in [Fig. 4,](#page-4-2) was produced by analysing the band of functional groups attached to glass using a deconvolution technique. The deconvolution peaks are related to the active peaks inside the glass network. The peaks can be coordinated as following: The combined bending vibration of Si-O-Si in [SiO<sub>4</sub>] causes a peak of about 480 cm<sup>-1</sup> [\[19](#page-9-9)–[21\]](#page-9-9). The band at 588-554  $\text{cm}^{-1}$ , which coincides with the vibrations of  $W-O$  bonds, relates to the O-Si-O mode. The Na<sup>+</sup> vibration within the glass network is what causes the band to form at 529  $cm^{-1}$ 

<span id="page-4-2"></span>

Fig. 4. Deconvoluted FTIR spectrum for S1 sample.



Fig. 3. The FTIR spectra of the glass samples.

<span id="page-5-0"></span>Table 2. FTIR absorption spectra analysis of the prepared glass samples.

	S1	S2	S3	S4	S5
Centre	490.00	480.49	474.00	479.07	469.77
Area	2.52	0.84	1.36	1.34	0.39
Centre		568.49	590.00	545.77	555.90
Area		2.01	2.32	2.57	2.72
Centre	685.00	704.43	707.43	710.00	720.00
Area	9.50	9.25	10.14	11.25	11.91
Centre	845.00	847.66	853.66	854.02	863.50
Area	16.61	13.33	10.51	9.20	8.90
Centre	950.00	950.37	958.37	960.00	965.00
Area	18.32	14.73	12.25	11.96	10.41
Centre	1020.00	1044.04	1051.04	1090.46	1020.43
Area	10.59	10.87	9.98	9.72	9.33
Centre	1133.00	1123.02	1128.02	1120.00	1114.35
Area	10.01	10.08	10.76	10.98	11.76
Centre	1280.00	1273.46	1262.46	1260.00	1258.00
Area	7.71	10.06	13.49	15.45	16.59
Centre	1463.00	1487.04	1481.04	1480.00	1477.00
Area	11.02	11.31	13.69	13.70	14.08
Centre	1683.00	1680.00	1666.00	1630.00	1627.13
Area	8.85	7.23	6.18	5.03	2.03
$\rm N_4$	0.48	0.42	0.34	0.32	0.29

[\[22](#page-9-10)]. The peak vibrations around 670  $cm^{-1}$  and 720 cm<sup>-1</sup> are related to B-O-B vibrations, which come from di-borate  $B_4O_7$  (670 cm<sup>-1</sup>) and pentaborate  $B_5O_8$  (720 cm<sup>-1</sup>) [[23,](#page-9-11)[24](#page-9-12)].

The band at 860  $cm^{-1}$  could be attributed to the stretching vibrations  $B-O$  in tetrahedral  $BO<sub>4</sub>$  in penta-borate groups [[25](#page-9-13)[,26](#page-9-14)]. The peak at 920  $cm^{-1}$ represents the  $Si-O^-$  stretching vibration of  $[SiO_4]$ tetrahedron, the  $Si-O-B$  stretching vibration and the  $B-O^-$  stretching vibration mode of  $[BO_4]$  tetra-hedron [[27,](#page-9-15)[28](#page-9-16)]. The band at 1020 to 1088  $\text{cm}^{-1}$  refers to the B-O stretching vibrations of  $BO<sub>4</sub>$  in tri-borate groups [\[29](#page-10-0)]. Overlapping of the asymmetric stretching vibration of Si-O in the  $[SiO<sub>4</sub>]$  tetrahedron, and the reduction in peak strength means a reduction in the bridging oxygen [[30,](#page-10-1)[31](#page-10-2)]. 1182  $cm^{-1}$ band referred to the  $B - O$  stretching vibrations of  $BO<sub>3</sub>$  orthoborate groups [[32,](#page-10-3)[33](#page-10-4)]. 1261 cm<sup>-1</sup> band attributed to the stretching vibrations of  $B - O^-$  in BO<sub>3</sub> [[34\]](#page-10-5). The band located at ~1420 and 1490 cm<sup>-1</sup> may be due to  $B-O$  stretching vibration in BO units [\[35](#page-10-6)]. The band at 1640 cm<sup>-1</sup> is O-H vibration [\[29](#page-10-0),[33\]](#page-10-4).  $N_4 = A_4/(A_3 + A_4)$  was used to calculate the amount of quadruple boron,  $N_4$ .  $A_4$  and  $A_3$  represent the relative areas of the  $BO_4$  and  $BO_3$ , respectively. Deconvoluted FTIR spectra are shown [Fig. 4](#page-4-2) and in details in [Table 2.](#page-5-0) The behaviour of N4 within the addition of  $WO<sub>3</sub>$  is shown in [Fig. 5.](#page-5-1)

#### 3.3. Optical measurements results

The band gap of optical energy  $(E_{\text{opt}})$ , the refractive index (n), and the Urbach energy  $(E_u)$  The optical parameters of the glass system, such as the optical band gap energy ( $E_{\text{opt}}$ ), Urbach energy ( $E_{\text{U}}$ ), and refractive index (n), are determined with the use of UV-Vis spectrophotometer measurements. Using the Beer-Lambert law  $[36]$  $[36]$ , the following features may be utilized to describe the structure of the glass host at any given light incident:

$$
I_t = I_0 e^{-\alpha t} \tag{8}
$$

(t) is the sample's thickness,  $(I_t)$  and  $(I_0)$  denote the light's incidence and transmittion intensities, respectively.  $I_0$  is the total amount of light.

<span id="page-5-1"></span>

Fig. 5. Dependence of  $N_4$  on the  $WO_3$  content.

The optical absorption spectra obtained using UV-Vis spectroscopy are shown in Fig. $(6)$  it is confirmed the amorphous nature of the prepared glass samples since they showed clear absorption edges. The transmission spectra of glass samples are also shown in [Fig. \(7\)](#page-6-1). Due to nearly the same shape of the spectra with increasing  $WO<sub>3</sub>$  concentration, this reveals that all the prepared glass samples have interesting optical properties [\[37](#page-10-8)].

Meanwhile, the Davis and Mott relation [\[38](#page-10-9)] can be used to find the optical band gap energy  $(E_{opt})$  of the glass samples:

$$
\alpha h\nu = B\left(h\nu - E_{opt}\right)^n\tag{9}
$$

where  $h$   $\nu$  is incident photon energy, B is a constant,  $E_{opt}$  is the optical band gap energy, and  $\alpha$  is the

<span id="page-6-0"></span>

<span id="page-6-1"></span>Fig. 6. The optical absorbance spectra for all the glass samples.

absorption coefficient. With  $n = 2$  or 1/2 for permitted indirect and direct transition, the value of n refers to Eopt.

Urbach energy  $(E_U)$  is the length of the absorption edge's exponential tail which is determined by the following empirical relation [\[39](#page-10-10)],

$$
\alpha = \alpha_0 e^{\frac{h\nu}{c}u} \tag{10}
$$

where  $E_U$  is the width of the band tails of localised states, v is the radiation frequency, and  $\alpha_0$  is a constant.

The values of the indirect optical band gap energy were calculated using the relationship between  $(\alpha h\nu)^{0.5}$  and (hv) [Fig. \(8\)](#page-7-0) [[38](#page-10-9)–[41](#page-10-9)]. On the other hand, as shown in Fig. (9) The calculation of the Urbach as shown in [Fig. \(9\)](#page-7-1), The calculation of the Urbach energy values [\[39](#page-10-10),[41,](#page-10-11)[42](#page-10-12)] was made using the relation between  $\ln (\alpha)$  and (hv). As seen in [Fig. \(10\)](#page-8-0), structural variations of the glasses under study caused the energy gap and Urbach energy data to be vary dramatically. In the considered borate network, the generation of non-bridging oxygens increased the disorder state.

The refractive index (n), of any substance is an important physical property. It basically has control over the material's optical and electrical characteristics. Additionally, it is directly related to the domestic fields present in the substance and the ions and the electronic polarizability (p). The study of the connection between (n) and (p) for any semiconducting material is of consider a scientific interest because it has some of optoelectronic



Fig. 7. The optical transmittance spectra for all the glass samples.

<span id="page-7-0"></span>

Fig. 8. The change of  $(\alpha h\nu)^{0.5}$  with hv for the prepared glass samples.

applications and devices [[42\]](#page-10-12). The higher polarizability of  $W^{3+}$  cations in glass system that caused the increase of (n) [\[43](#page-10-13)].

The refractive index (n) calculated using the following formula using the optical band gap energy (Eopt) [\[44](#page-10-14)],

Molar refraction  $(R_m)$  is the result of reflection loss  $((n^2-1)/(n^2+2))$  and molar volume  $(V_m)$ . This parameter is produced by the Lorentz-Lorentz equation, where  $R_m$  is proportional to  $V_m$ , and it has the following implications for the glass structure [\[45](#page-10-15)],

<span id="page-7-1"></span>
$$
\frac{n^2 - 1}{n^2 + 2} = 1 - \sqrt{\frac{E_{opt}}{20}}\tag{12}
$$



Fig. 9. The variation of ln  $(\alpha)$  with hv of the prepared glass samples.

<span id="page-8-0"></span>

Fig. 10. The relation of  $E_{opt}$  and  $E_u$  with  $WO_3$  content.

The molar electronic polarizability am, according to Clausios Mosotti, is given by the relation:

$$
\alpha_m = \left(\frac{3}{4\pi N_A}\right) R_m \tag{13}
$$

where  $N_A$  is Avogadro's number.

Optical basicity  $(\Lambda)$  and electronic polarizability  $(\alpha_0^2)$  values can be obtained of Eg by using the formulas as: [[46](#page-10-16)[,47](#page-10-17)].

$$
\alpha_0^{2-} = -0.24192E_{opt} + 3.5\tag{14}
$$

$$
\Lambda = -0.1344E_{opt} + 1.7\tag{15}
$$

As the  $W^{3+}$  concentration was increased, the electronic polarizability of oxide and the optical basicity increase ([Table 3\)](#page-8-1). This indicates that the oxygen ions in the present glass system have the chance to transfer electrons to the surrounding cations [\[47](#page-10-17)].

<span id="page-8-1"></span>Table 3. Optical parameters in the glasses sample.

Optical parameter	S1	S <sub>2</sub>	S3	S4	S5
$E_{opt}$ (eV)	3.30	2.75	2.54	2.22	1.89
$E_{II}$ (eV)	1.70	2.48	2.72	3.23	4.82
Refractive index, n	1.98	2.10	2.15	2.24	2.35
optical electronegativity (γ)	0.89	0.74	0.68	0.60	0.51
Metallization	0.51	0.47	0.45	0.43	0.40
dielectric constant	3.92	4.41	4.62	5.01	5.52
electronic polarizability $(\alpha_0^2)$	2.70	2.83	2.89	2.96	3.04
optical basicity $(\Lambda)$	1.26	1.33	1.36	1.40	1.45

The optical electronegativity  $(\chi)$  can be calculated based on  $E_{opt}$  by applying this relation [\[47](#page-10-17)]:

$$
\chi = 0.2688 E_{opt} \tag{16}
$$

The kind of bonding existing in the materials is also indicated by optical electronegativity  $(\chi)$ . The values  $\chi$  are found to decrease when WO<sub>3</sub> levels increases in the network glasses. This decrease is brought on by its inverse relationship to the refractive index. The modest changes in electronegativity values might perhaps be caused by the system's covalent nature [\[48](#page-10-18),[49\]](#page-10-19).

Metallization criterion (M) gives the non-metallic or metallic nature of the oxide glasses, where Rm/  $Vm > 1$  for metal and <1 for non-metal [[49\]](#page-10-19). Higher values of M are the trend of material or compound that has larger optical band gap  $(E_{opt})$  and smaller refractive index (n) and vice versa [[50](#page-10-20)[,51](#page-10-21)]. Insulating material has M near 1. A relation can be derived from the Lorentz-Lorenz equation,

$$
M = 1 - \frac{R_m}{V_m} \tag{17}
$$

where  $V_m$  is the molar volume of the respective glass sample. The M data showed almost linear decrease as  $WO<sub>3</sub>$  increased, these results are shown in [Table 3.](#page-8-1) However, the increment in  $WO<sub>3</sub>$  concentration in the glass system caused decrease in M values due to the contraction of  $E_{opt}$  [\[47](#page-10-17)].

Optical dielectric constant,  $\mathcal{E}_{opt}$  can be calculated using the equation:

$$
\epsilon_{opt} = n^2 \tag{18}
$$

According to [Table 3](#page-8-1), the  $\mathcal{E}_{opt}$  showed increasing trend with the  $WO_3$  increasing from 3.91 to 6.50. The increase may be due to the rise in NBO in the network of glass samples due to the addition of  $WO<sub>3</sub>$ .

#### 4. Conclusion

From this work we can conclude that the FTIR result showed that  $BO_3$  increased and  $BO_4$ decreased with increasing  $WO<sub>3</sub>$  content, accordingly, NBO increases. The density measurements confirmed the glassy state of the prepared samples. The optical band gap, optical electronegativity and Metallization showed decreased values while Urbach energy, Optical band gap show high value, 3.3 eV for  $x = 0$  which decreases to 1.89 eV for sample 8 mol % with increase in refractive index from 1.98 up to 2.35 the refractive index, dielectric constant, electronic polarizability and optical basicity increased with increasing  $WO<sub>3</sub>$  content. These results let us conclude that these samples are promising in electrical, energy storage and optoelectronic applications.

#### Conflicts of interest

The authors have no conflict of interest to declare related to the content of the article.

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